

$[Q, R] = 0$ and Verlinde Series

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Joint with Eckhard Meinrenken

The Hamiltonian $[Q, R] = 0$ Theorem

- G compact Lie group, Lie algebra \mathfrak{g} , representation ring $R(G)$
- $\phi : M \rightarrow \mathfrak{g}^*$ compact Hamiltonian G -space
- $L \rightarrow M$ prequantum line bundle, G -equivariant
- D_L Dolbeault-Dirac operator for compatible almost \mathbb{C} structure, twisted by L
- $Q(M) \in R(G)$ the index of D_L

Theorem (Meinrenken) “Guillemin-Sternberg principle”

Assume 0 is a regular value. Then

$$Q(M)^G = Q(M//G) \quad (\text{where } M//G = \phi^{-1}(0)/G)$$

- Paradan (2001) decomposed the index $Q(M) \in R(G)$ into a sum of simpler, infinite-dimensional contributions $Q_\beta(M) \in R^{-\infty}(G)$ indexed by the components of the critical set of $\|\phi\|^2$.
- Szenes-Vergne (2010) re-derived Paradan's decomposition using a combinatorial rearrangement of the Atiyah-Bott fixed-point formula, together with asymptotics of $Q(M, L^k)$, $k \rightarrow \infty$.

$$M = \mathbf{CP}^1 \circlearrowleft S^1 = G, t \cdot [x : y] = [tx : t^{-1}y], k > 0$$

$$Q(M, L^k)(t) = \frac{t^k}{1 - t^{-2}} + \frac{t^{-k}}{1 - t^2} = \sum_{n=0}^k t^{-k+2n}.$$

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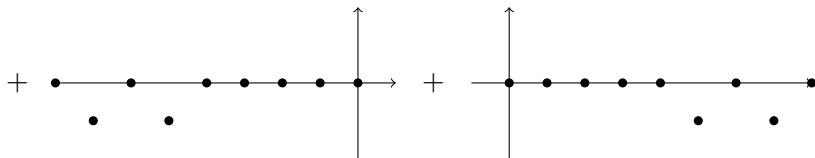
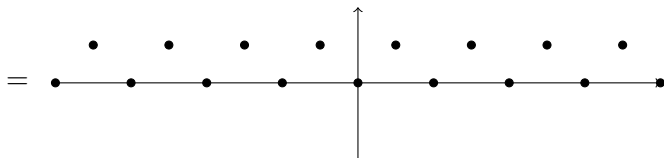
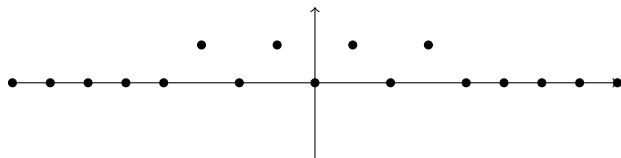
Paradan's formula in terms of multiplicities:

$$\begin{aligned} m(\lambda, k) &= \begin{cases} 1 & |\lambda| \leq k, \lambda = k \pmod{2} \\ 0 & \text{else} \end{cases} \\ &= m_{\text{q-pol}}(\lambda, k) + m_+(\lambda, k) + m_-(\lambda, k) \end{aligned}$$

m_{\pm} are corrections, supported away from $0 \in \mathbb{Z} = \Lambda^*$.

Example

$k = 3$



Let Γ be a lattice.

Definition

A function $f : \Gamma \rightarrow \mathbb{C}$ is **quasi-polynomial** if there is a full-rank sublattice $\Gamma' \subset \Gamma$ such that f is polynomial on each coset $\gamma + \Gamma'$.

Example

$$m_{\text{q-pol}}(\lambda, k) = \begin{cases} 1 & \lambda = k \pmod{2} \\ 0 & \text{else} \end{cases}$$

is quasi-polynomial on $\Lambda^* \times \mathbb{Z}$.

Important! Study multiplicities as function of $(\lambda, k) \in \Lambda^* \times \mathbb{N}$.

Weyl character formula

$$\chi_\lambda|_{\mathcal{T}} = \frac{\sum_w (-1)^{\ell(w)} e_{w(\lambda+\rho)-\rho}}{\prod_{\alpha>0} 1 - e^{-\alpha}}, \quad \lambda \in \Lambda^* \cap \mathfrak{t}_+^*$$

Multiplicity function m for $Q(M) \in R(G)$

$$m : \Lambda^* \cap \mathfrak{t}_+^* \rightarrow \mathbb{Z}$$

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Instead study

$$\prod_{\alpha>0} (1 - e_{-\alpha}) Q(M)|_T \in R(T)$$

→ multiplicity function m^a is W -anti-symmetric (shifted action),
and

$$m^a(\lambda) = m(\lambda), \quad \lambda \in \Lambda^* \cap \mathfrak{t}_+^*$$

$\phi : M \rightarrow \mathfrak{g}^*$ transverse to $\mathfrak{t}^* \Rightarrow$

$$\prod_{\alpha > 0} (1 - e_{-\alpha}) Q(M)|_{\mathcal{T}} = Q(X), \quad X = \phi^{-1}(\mathfrak{t}^*)$$

Hamiltonian LG -spaces

- $LG = \text{Map}(S^1, G)$ loop group, $L\mathfrak{g} = \Omega^0(S^1, \mathfrak{g})$
- $L\mathfrak{g}^* = \Omega^1(S^1, \mathfrak{g})$, LG acts by gauge transformations:

$$g \cdot \xi = \text{Ad}_g \xi - dg g^{-1}.$$

Definition

A *Hamiltonian LG -space* $(\mathcal{M}, \omega, \Psi)$ consists of a symplectic (Banach) LG -manifold, equipped with a proper moment map $\Psi : \mathcal{M} \rightarrow L\mathfrak{g}^*$.

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Example: Moduli of flat connections

Σ compact Riemann surface with b boundary components \Rightarrow

$$\mathcal{M} = \mathcal{A}^{\text{flat}}(\Sigma) / \{g \in \mathcal{G} : g|_{\partial\Sigma} = 1\}, \quad \Psi(A) = \iota_{\partial\Sigma}^* A$$

is a Hamiltonian LG^b -space.

Let $L \rightarrow \mathcal{M}$ be a prequantum line bundle (\widehat{LG} -equivariant).

\exists “quantization” method (Meinrenken '10) for the finite-dimensional quotient:

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(M no longer symplectic, instead **quasi-Hamiltonian** (AMM '98))

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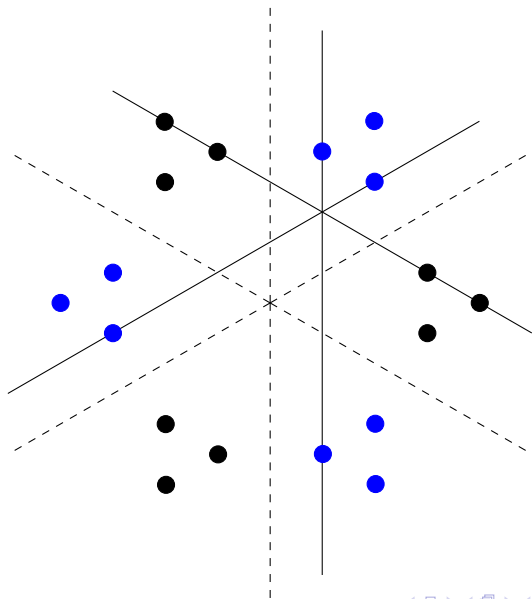
The analog of m^a exists... It is a function:

$$m^a : \Lambda^* \rightarrow \mathbb{Z}$$

which is **W_{aff}-anti-symmetric** (shifted). Take powers $L^k \Rightarrow$

$$m^a : \Lambda^* \times \mathbb{N} \rightarrow \mathbb{Z}$$

m^a for a Hamiltonian $LSU(3)$ -space



Formula for m^a

$$m^a(\lambda, k) = \frac{1}{|T_\ell|} \sum_{t \in T_\ell} t^{-\lambda} \prod_{\alpha > 0} (1 - t^{-\alpha}) Q(M, k)(t)$$

$$Q(M, k)(t) = \sum_{F \subset M} \int_F \frac{\hat{A}(F) \text{Ch}(\mathcal{L}_k, t)^{1/2}}{\mathcal{D}_{\mathbb{R}}(\nu_F, t)}$$

- $\ell = k + h^\vee$, h^\vee dual Coxeter number (e.g. $h^\vee = N$ for $SU(N)$)
- T_ℓ finite “Verlinde” subgroup

$$T_\ell := \ell^{-1} B^\sharp(\Lambda^*) / \Lambda \subset \mathfrak{t} / \Lambda = T$$

Goal:

Prove $m^a(0, k)$ is **quasi-polynomial**, using strategy inspired by Szenes-Vergne-Paradan. Derive $[Q, R] = 0$ from this + stationary phase argument for large k .

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Challenges:

- Finite group $T_\ell \subset T$, varying with ℓ .
- $\|\Psi\|^2$ on \mathcal{M} does not descend to M .

$\Lambda \subset \Xi$ lattices in vector space \mathfrak{h}

$\alpha = (\alpha_1, \dots, \alpha_n)$, $\alpha_k \in \Lambda^*$.

Definition

The *Verlinde series* $V(\alpha) : \Lambda^* \times \mathbb{N} \rightarrow \mathbb{C}$ is

$$V(\alpha)(\lambda, \ell) = \sum'_{\xi \in \ell^{-1}\Xi/\Lambda} \frac{e^{-2\pi i \langle \lambda, \xi \rangle}}{\prod_{\alpha} 1 - e^{2\pi i \langle \alpha, \xi \rangle}}.$$

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Szenes' ('03): remarkable residue formula for $V(\alpha)$. Implies $V(\alpha)$ is **piecewise quasi-polynomial**.

\Rightarrow for generic γ , $V(\alpha)$ has a **quasi-polynomial “germ”**
 $\text{Ver}(\alpha, \gamma)$.

Expansion in direction +1:

$$\frac{1}{1-t} = \sum_{n \geq 0} t^n$$

Expansion in direction -1:

$$\frac{1}{1-t} = \frac{-t^{-1}}{1-t^{-1}} = -t^{-1} \sum_{n \geq 0} t^{-n}$$

A combinatorial decomposition

$\alpha = (\alpha_1, \dots, \alpha_n)$. Choose generic vector γ^+ , and define $P(\alpha, \gamma^+)(\lambda)$ to be the coefficient of t^λ in expansion of

$$\frac{1}{\prod_k (1 - t^{\alpha_k})}$$

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Theorem (Boysal-Vergne)

Let \mathcal{S} be the set of affine subspaces generated by Ξ^* , α . For generic γ ,

$$V(\alpha) = \sum_{\Delta \in \mathcal{S}} \text{Ver}(\alpha_\Delta, \gamma_\Delta) * P(\alpha'_\Delta, \gamma_\Delta^+)$$

where $\gamma_\Delta = \text{proj}_\Delta(\gamma)$, $\gamma_\Delta^+ = \gamma_\Delta - \gamma$ and $\alpha_\Delta \cup \alpha'_\Delta = \alpha$.

Recall $T_\ell = \ell^{-1} B^\sharp(\Lambda^*)/\Lambda$, and

$$m^a(\lambda, k) = \frac{1}{|T_\ell|} \sum_{t \in T_\ell} t^{-\lambda} \prod_{\alpha > 0} (1 - t^{-\alpha}) \sum_{F \subset M^t} \int_F \frac{\hat{A}(F) \text{Ch}(\mathcal{L}_k, t)^{1/2}}{\mathcal{D}_{\mathbb{R}}(\nu_F, t)}.$$

Lemma

*The multiplicity m^a is a sum of a (possibly large number) of **translates** of Verlinde series, for various infinitesimal stabilizers $\mathfrak{h} \subset \mathfrak{t}$, weighted by characteristic numbers of the fixed-point sets.*

Rearrangement of fixed-point formula I

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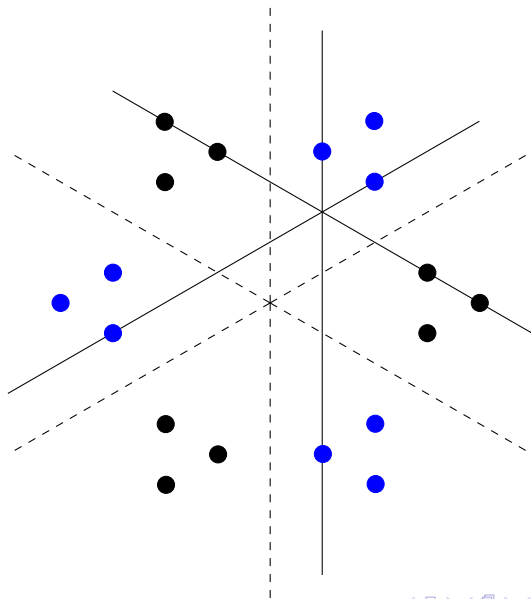
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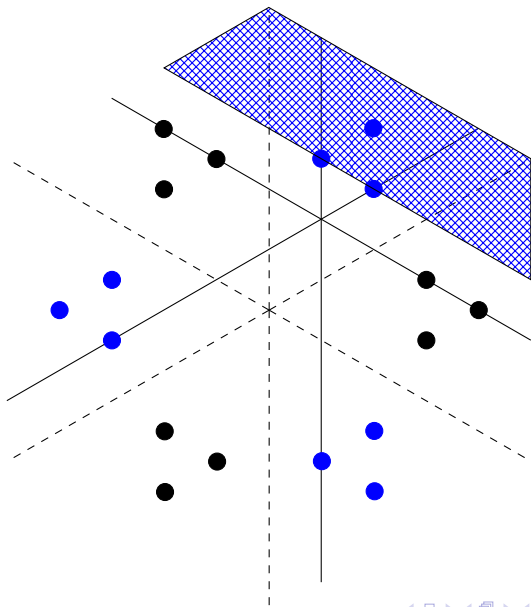
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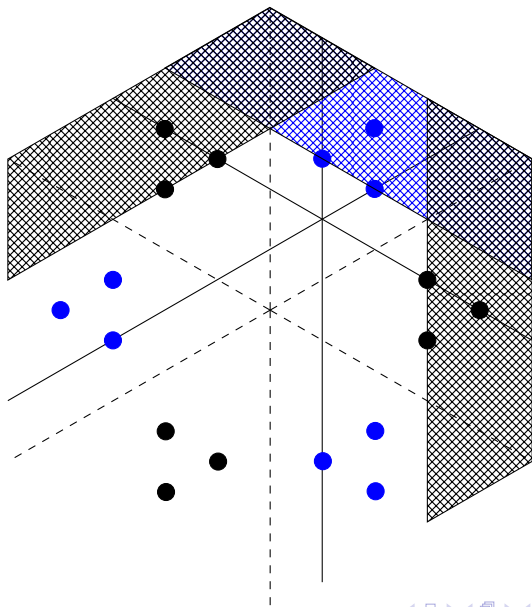
Apply Boysal-Vergne formula term-by-term \rightarrow complicated formula for m^a .

m^a for a Hamiltonian $LSU(3)$ -space



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Schematically:

$$m^a = m_{\text{q-pol}}^a + \sum_{\Delta} m_{\Delta}^a,$$

where $\Delta \in \mathcal{S}(M)$ a collection of proper affine subspaces of \mathfrak{t}^* .

Each term m_{Δ}^a involves a further sum over fixed-point sets $F \subset M^{\mathfrak{t}_{\Delta}}$, $\mathfrak{t}_{\Delta} = \Delta^{\perp}$.

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Moreover:

- central contribution $m_{\text{q-pol}}^a$ is quasi-polynomial
- each $m_{\Delta}^a(-, k)$ supported in a half-space
- for $k \gg 0$, these half-spaces do not contain 0

Quasi-polynomial?

$$m^a = m_{\text{q-pol}}^a + \sum_{\Delta} m_{\Delta}^a$$

If we show

$$0 \notin \text{supp}(m_{\Delta}^a) \quad \text{for all } k \geq 1$$

then

$$m^a(0, k) = m_{\text{q-pol}}^a(0, k) \quad \forall k \geq 1 \quad \Rightarrow \quad m^a \text{ quasi-polynomial!}$$

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- Bound translations (“rho shifts”) occurring above.
- Prove remaining problematic terms (those not corresponding to critical points of $\|\Psi\|^2$) **cancel out**.

The Stiefel diagram is an affine hyperplane arrangement in \mathfrak{t} :

$$\langle \alpha, x \rangle + n = 0, \quad \alpha \in \mathcal{R}, n \in \mathbb{Z}.$$

v a vertex and $K = G_{\exp(v)}$. Let ρ, ρ_K be half-sums of positive roots for G, K , and let

$$c_K = \frac{\rho + \rho_K}{2h^v}.$$

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Theorem: Identify $\mathfrak{t} \simeq \mathfrak{t}^*$ (basic inner product). Then

$$\|v - c_K\| \geq \|c_K\|.$$

Special case: $G = SU(N)$, then $c_K = c$ is the barycenter of the alcove.

$$m^a = m_{q\text{-pol}}^a + \sum_{\Delta} m_{\Delta}^a \quad (1)$$

Back to $\Psi : \mathcal{M} \rightarrow Lg^*$

$$\text{Crit}(\|\Psi\|^2) = G \cdot \bigcup_{\beta \in \mathcal{B}} \mathcal{M}^{\beta} \cap \Psi^{-1}(\beta).$$

$\beta = 0$ corresponds to $m_{q\text{-pol}}$, and $\beta \in \Delta$ corresponds to m_{Δ} .

A priori, there are **additional problematic terms** in (1) that do not correspond to any $\beta \in \mathcal{B}$.

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These terms vanish! (Involves using fixed-point formula in reverse on submanifolds $C \subset M^{\Delta}$, and asymptotic argument $k \rightarrow \infty$.)

Asymptotic $[Q, R] = 0$ (rough idea)

Atiyah-Segal-Singer fixed-point formula \rightarrow Berline-Vergne
“delocalized” index formulas, and stationary phase $k \rightarrow \infty$
(Meinrenken '94)

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Recall $\Psi : \mathcal{M} \rightarrow Lg^*$. Assume $X = \Psi^{-1}(\mathfrak{t}^*)$ smooth. For $\xi \in \mathfrak{t}$
small, consider

$$\chi(\xi) := \int_X \hat{A}(X, \xi) \text{Ch}(\mathcal{L}_k, \xi)$$

Integral lattice Λ acts on X . Integrand periodic except for phase shift:

$$\chi(\xi) = \sum_{\lambda} e^{2\pi i \ell \langle \lambda, \xi \rangle} \int_{X_0} \hat{A}(X, \xi) \text{Ch}(\mathcal{L}_k, \xi)$$

where X_0 is a fundamental domain.

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Poisson summation formula $\Rightarrow \chi$ supported on $\ell^{-1}B^{\sharp}(\Lambda^*)$.

Recall

$$T_{\ell} = \ell^{-1}B^{\sharp}(\Lambda^*)/\Lambda$$

Localize on $X/\Lambda \Rightarrow$ fixed-point formulas for $Q(M, k)(t)$,
 $t = \exp(\xi)$.

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Stationary phase argument \Rightarrow

$$\int_t \chi(\xi) d\xi = \int_t d\xi \int_X \hat{A}(X, \xi) \text{Ch}(\mathcal{L}_k, \xi)$$

concentrated near $\Psi^{-1}(0)$ as $k \rightarrow \infty$.